

Übungen zur Theoretischen Mechanik
 FSU Jena - SS 2007
 Blatt 02 - Lösungen

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(2) Massenpunkt auf Kreis

In Kugelkoordinaten gilt im allgemeinen:

$$\vec{r} = r \cdot \begin{bmatrix} \sin \theta \cos \varphi \\ \sin \theta \sin \varphi \\ \cos \theta \end{bmatrix}$$

$$\vec{g}_r = \frac{\partial \vec{r}}{\partial r} = \begin{bmatrix} \sin \theta \cos \varphi \\ \sin \theta \sin \varphi \\ \cos \theta \end{bmatrix}, \quad \vec{g}_\theta = \frac{\partial \vec{r}}{\partial \theta} = r \cdot \begin{bmatrix} \cos \theta \cos \varphi \\ \cos \theta \sin \varphi \\ -\sin \theta \end{bmatrix}, \quad \vec{g}_\varphi = \frac{\partial \vec{r}}{\partial \varphi} = r \cdot \begin{bmatrix} -\sin \theta \sin \varphi \\ \sin \theta \cos \varphi \\ 0 \end{bmatrix}$$

$$\lambda_r := |\vec{g}_r| = 1, \quad \lambda_\theta := |\vec{g}_\theta| = r, \quad \lambda_\varphi := |\vec{g}_\varphi| = r \sin \theta$$

$$\vec{e}_r := \frac{\vec{g}_r}{\lambda_r}, \quad \vec{e}_\theta := \frac{\vec{g}_\theta}{\lambda_\theta}, \quad \vec{e}_\varphi := \frac{\vec{g}_\varphi}{\lambda_\varphi}$$

$$\dot{\vec{e}}_r = \begin{bmatrix} \dot{\theta} \cos \theta \cos \varphi - \dot{\varphi} \sin \theta \sin \varphi \\ \dot{\theta} \cos \theta \sin \varphi + \dot{\varphi} \sin \theta \cos \varphi \\ -\dot{\theta} \sin \theta \end{bmatrix} = \dot{\theta} \cdot \vec{e}_\theta + \dot{\varphi} \sin \theta \cdot \vec{e}_\varphi$$

$$\dot{\vec{e}}_\theta = \begin{bmatrix} -\dot{\theta} \sin \theta \cos \varphi - \dot{\varphi} \cos \theta \sin \varphi \\ -\dot{\theta} \sin \theta \sin \varphi + \dot{\varphi} \cos \theta \cos \varphi \\ -\dot{\theta} \cos \theta \end{bmatrix} = -\dot{\theta} \cdot \vec{e}_r + \dot{\varphi} \cos \theta \cdot \vec{e}_\varphi$$

$$\dot{\vec{e}}_\varphi = \begin{bmatrix} -\dot{\varphi} \cos \varphi \\ -\dot{\varphi} \sin \varphi \\ 0 \end{bmatrix} = -\dot{\varphi} \sin \theta \cdot \vec{e}_r - \dot{\varphi} \cos \theta \cdot \vec{e}_\theta$$

$$\dot{\vec{r}} = \dot{r} \cdot \vec{g}_r + \dot{\theta} \cdot \vec{g}_\theta + \dot{\varphi} \cdot \vec{g}_\varphi = \dot{r} \cdot \vec{e}_r + r\dot{\theta} \cdot \vec{e}_\theta + r\dot{\varphi} \sin \theta \cdot \vec{e}_\varphi$$

$$\ddot{\vec{r}} = \left(\ddot{r} - r\dot{\theta}^2 - r\dot{\varphi}^2 \sin^2 \theta \right) \cdot \vec{e}_r + \left(r\ddot{\theta} + 2\dot{r}\dot{\theta} - r\dot{\varphi}^2 \sin \theta \cos \theta \right) \cdot \vec{e}_\theta + \left(r\ddot{\varphi} \sin \theta + 2\dot{r}\dot{\varphi} \sin \theta + 2r\dot{\theta}\dot{\varphi} \cos \theta \right) \cdot \vec{e}_\varphi$$

• **Kugelkoordinaten:**

$$r = R_0 : \text{const} \Rightarrow \dot{r} = 0, \quad \dot{\phi} = \omega : \text{const} \Rightarrow \ddot{\phi} = 0$$

$$\Rightarrow \dot{\vec{r}} = R_0 \dot{\theta} \cdot \vec{e}_\theta + \omega R_0 \sin \theta \cdot \vec{e}_\varphi \Rightarrow |\dot{\vec{r}}| = v_0 = R_0 \sqrt{\dot{\theta}^2 + \omega^2 \sin^2 \theta} : \text{const} \Rightarrow \dot{\theta} = \sqrt{\frac{v_0^2}{R_0^2} - \omega^2 \sin^2 \theta} = \omega \cos \theta \quad (\text{da } v_0 = \omega R_0)$$

$$\ddot{\theta} = \frac{d(\omega \cos \theta)}{dt} = -\omega \dot{\theta} \sin \theta = -\omega^2 \cos \theta \sin \theta$$

$$\Rightarrow \dot{\vec{r}} = v_0 \cos \theta \cdot \vec{e}_\theta + v_0 \sin \theta \cdot \vec{e}_\varphi$$

$$\begin{aligned} \ddot{\vec{r}} &= -(\omega^2 R_0 \cos^2 \theta + \omega^2 R_0 \sin^2 \theta) \cdot \vec{e}_r - (\omega^2 R_0 \cos \theta \sin \theta + \omega^2 R_0 \sin \theta \cos \theta) \cdot \vec{e}_\theta + 2\omega^2 R_0 \cos^2 \theta \cdot \vec{e}_\varphi \\ &= \omega^2 R_0 \cdot (-\vec{e}_r - 2 \cos \theta \sin \theta \cdot \vec{e}_\theta + 2 \cos^2 \theta \cdot \vec{e}_\varphi) \end{aligned}$$

Normalkoordinaten:

Sind \vec{T} der Bahn-Tangenteneinheitsvektor und \vec{N} der Bahn-Normaleinheitsvektor, so gilt

$$\dot{\vec{r}} = v_0 \cdot \vec{T} \Rightarrow \vec{T} = \cos \theta \cdot \vec{e}_\theta + \sin \theta \cdot \vec{e}_\varphi = \begin{pmatrix} \cos^2 \theta \cos \varphi - \sin \theta \sin \varphi \\ \cos^2 \theta \sin \varphi + \sin \theta \cos \varphi \\ -\cos \theta \sin \theta \end{pmatrix}$$

$$\ddot{\vec{r}} = v_0 \cdot \vec{T} + \frac{v_0^2}{R} \cdot \vec{N} = \frac{v_0^2}{R} \cdot \vec{N} \Rightarrow \frac{v_0^2}{R} = |\ddot{\vec{r}}| \Rightarrow R = \frac{v_0^2}{|\ddot{\vec{r}}|} = \frac{R_0}{\sqrt{1 + 4 \cos^2 \theta}}$$

$$\Rightarrow \vec{N} = \frac{-\vec{e}_r - 2 \cos \theta \sin \theta \cdot \vec{e}_\theta + 2 \cos^2 \theta \cdot \vec{e}_\varphi}{\sqrt{1 + 4 \cos^2 \theta}} = \frac{1}{\sqrt{1 + 4 \cos^2 \theta}} \cdot \begin{bmatrix} -\sin \theta \cos \varphi (1 + 2 \cos^2 \theta) - 2 \cos^2 \theta \sin \varphi \\ -\sin \theta \sin \varphi (1 + 2 \cos^2 \theta) + 2 \cos^2 \theta \cos \varphi \\ \cos \theta (2 \sin^2 \theta - 1) \end{bmatrix}$$

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$$\dot{\theta} = \omega \cos \theta \Rightarrow \int_{\theta_0}^{\theta} \frac{du}{\cos u} = \int_{t_0}^t \omega d\lambda \Rightarrow \ln \left| \tan \left(\frac{\theta}{2} + \frac{\pi}{4} \right) \right| = \omega(t - t_0) + C, \quad C \in \mathbb{R}$$

$$\Rightarrow \theta(t) = 2 \arctan \left(A e^{\omega(t-t_0)} \right) - \frac{\pi}{2}, \quad A \in \mathbb{R}^+$$

$$R(t) = \frac{R_0}{\sqrt{1 + 4 \cos^2 \theta(t)}} = \frac{R_0}{\sqrt{1 + \frac{4A^2 e^{2\omega(t-t_0)}}{(1+A^2 e^{2\omega(t-t_0)})^2}}}$$

$$\text{Für } (\theta_0, t_0) \stackrel{!}{=} (0, 0) \rightarrow A = 1 \rightarrow \theta(t) = 2 \arctan(e^{\omega t}) - \frac{\pi}{2}, \quad R(t) = \frac{R_0}{\sqrt{1 + \frac{4e^{2\omega t}}{(1+e^{2\omega t})^2}}}$$

$$\lim_{t \rightarrow \infty} \theta = \frac{\pi}{2}, \quad \lim_{t \rightarrow \infty} R = \lim_{\theta \rightarrow \pi/2} \frac{R_0}{\sqrt{1 + 4 \cos^2 \theta}} = R_0$$

(3) Harmonische Schwingungen

$$f(t) = \alpha_0 |t| = \begin{cases} \alpha_0 t & : t \geq 0 \\ -\alpha_0 t & : t < 0 \end{cases} = \sum_{n=-\infty}^{\infty} A_n e^{in\omega t}, \quad \omega := \frac{2\pi}{T}$$

$$\text{Für } n \neq 0 : A_n = \frac{1}{T} \cdot \int_{-T/2}^{T/2} f(t) e^{-in\omega t} \cdot dt = \frac{\alpha_0}{T} \cdot \left[-\int_{-T/2}^0 t e^{-in\omega t} \cdot dt + \int_0^{T/2} t e^{-in\omega t} \cdot dt \right]$$

$$= \frac{\alpha_0}{T} \cdot \left[\int_0^{-T/2} t e^{-in\omega t} \cdot dt + \int_0^{T/2} t e^{-in\omega t} \cdot dt \right]$$

$$= \frac{\alpha_0}{T} \cdot \left[e^{-in\omega t} \cdot \left(\frac{it}{\omega n} + \frac{1}{n^2 \omega^2} \right) \right]_0^{-\pi/\omega} + \frac{\alpha_0}{T} \cdot \left[e^{-in\omega t} \cdot \left(\frac{it}{\omega n} + \frac{1}{n^2 \omega^2} \right) \right]_0^{\pi/\omega}$$

$$= \frac{2\alpha_0}{T} \cdot \left[\frac{(-1)^n}{n^2 \omega^2} - \frac{1}{n^2 \omega^2} \right] = \begin{cases} 0 & : n \text{ gerade} \\ -\frac{2\alpha_0 T}{n^2 \pi^2} & : n \text{ ungerade} \end{cases}$$

$$A_0 = \frac{\alpha_0}{T} \cdot \left[\int_0^{-T/2} t \cdot dt + \int_0^{T/2} t \cdot dt \right] = \frac{\alpha_0}{T} \cdot \left(\left[\frac{t^2}{2} \right]_0^{-T/2} + \left[\frac{t^2}{2} \right]_0^{T/2} \right) = \frac{\alpha_0 T}{4}$$

Es gilt: $A_n = A_{-n} \quad \forall n \in \mathbb{N}$

Demzufolge:

$$f(t) = \sum_{n=-\infty}^{\infty} A_n e^{in\omega t} = A_0 + \sum_{n=1}^{\infty} (A_n e^{in\omega t} + A_{-n} e^{-in\omega t}) = A_0 + \sum_{n=1}^{\infty} 2A_n \cos(n\omega t)$$

$$= \frac{\alpha_0 T}{4} - \frac{2\alpha_0 T}{\pi^2} \cdot \sum_{n=1}^{\infty} \frac{\cos((2n-1)\omega t)}{(2n-1)^2}$$