

Quantenmechanik II

FSU Jena - WS 2009/2010

Übungsserie 06 - Lösungen

Stilianos Louca

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Aufgabe 12

Mit

$$\mathbf{B} := \nabla \times \mathbf{A} \quad , \quad \mathbf{E} := -\nabla \Phi - \dot{\mathbf{A}}$$

als tatsächliche physikalische Felder, ist ersichtlich dass für beliebiges Skalarfeld $U(t, \mathbf{x})$, $U \in \mathcal{C}^2$ auch

$$\mathbf{A}' := \mathbf{A} + \nabla U \quad , \quad \Phi' := \Phi - \dot{U}$$

die gleichen $\mathbf{B}$, $\mathbf{E}$ -Felder definieren, denn

$$\nabla \times \mathbf{A}' = \nabla \times \mathbf{A} + \underbrace{\nabla \times (\nabla U)}_0 = \mathbf{B}$$

$$-\nabla \Phi' - \dot{\mathbf{A}}' = -\nabla \Phi - \dot{\mathbf{A}} + \cancel{\nabla \partial_t U} - \cancel{\partial_t \nabla U}$$

Gesucht ist nun ein geeignetes U mit $\nabla \cdot \mathbf{A}' = 0 = \Phi'$. Wegen

$$\partial_t \left[\nabla \cdot \mathbf{A}(t, \mathbf{x}) + \int_0^t \Delta \Phi(\tau, \mathbf{x}) \, d\tau \right] = -\nabla \cdot \underbrace{\partial_t \left[-\mathbf{A}(t, \mathbf{x}) - \int_0^t \nabla \Phi(\tau, \mathbf{x}) \, d\tau \right]}_{\mathbf{E}(t, \mathbf{x})} \stackrel{\text{Vakuum}}{=} 0$$

ist $\nabla \cdot \mathbf{A}(t, \mathbf{x}) + \int_0^t \Delta \Phi(\tau, \mathbf{x}) \, d\tau$ Zeitunabhängig. Dabei ist die Poissongleichung

$$\Delta h(\mathbf{x}) = - \left[\nabla \cdot \mathbf{A}(t, \mathbf{x}) + \int_0^t \Delta \Phi(\tau, \mathbf{x}) \, d\tau \right] \tag{0.1}$$

(bei hinreichend glatten Feldern) stets lösbar und der Ansatz

$$U(t, \mathbf{x}) := \int_0^t \Phi(\tau, \mathbf{x}) \, d\tau + h(\mathbf{x}) \tag{0.2}$$

erfüllt genau die obigen Forderungen, denn

$$\nabla \cdot \mathbf{A}' = \nabla \cdot \mathbf{A} + \Delta U \stackrel{(0.2)}{=} \nabla \cdot \mathbf{A} + \int \Delta \Phi \, dt + \Delta h \stackrel{(0.1)}{=} 0$$

$$\Phi' = \Phi - \partial_t U \stackrel{(0.2)}{=} 0$$

□

Aufgabe 13

Beginnen mit der Definition

$$|n_1, \dots, n_{\boldsymbol{\mu}}, \dots\rangle := \dots \frac{(\hat{a}_{\boldsymbol{\mu}}^\dagger)^{n_{\boldsymbol{\mu}}}}{\sqrt{n_{\boldsymbol{\mu}}!}} \dots \frac{(\hat{a}_1^\dagger)^{n_1}}{\sqrt{n_1!}} |0\rangle$$

und

$$\hat{\mathbf{E}}(t, \mathbf{x}) := -\partial_t \hat{\mathbf{A}}(t, \mathbf{x}) = \sum_{\boldsymbol{\mu}} \left(\frac{\hbar \omega_{\boldsymbol{\mu}}}{2V \epsilon_0} \right)^{\frac{1}{2}} \mathbf{e}_{\boldsymbol{\mu}} \cdot i \cdot [e^{i\mathbf{k}_{\boldsymbol{\mu}} \mathbf{x} - i\omega_{\boldsymbol{\mu}} t} \hat{a}_{\boldsymbol{\mu}} - e^{-i\mathbf{k}_{\boldsymbol{\mu}} \mathbf{x} + i\omega_{\boldsymbol{\mu}} t} \hat{a}_{\boldsymbol{\mu}}^\dagger]$$

$$\hat{\mathbf{B}}(t, \mathbf{x}) := \nabla \times \hat{\mathbf{A}}(t, \mathbf{x}) = \sum_{\boldsymbol{\mu}} \left(\frac{\hbar}{2V \epsilon_0 \omega_{\boldsymbol{\mu}}} \right)^{\frac{1}{2}} (\mathbf{k}_{\boldsymbol{\mu}} \times \mathbf{e}_{\boldsymbol{\mu}}) \cdot i \cdot [e^{i\mathbf{k}_{\boldsymbol{\mu}} \mathbf{x} - i\omega_{\boldsymbol{\mu}} t} \hat{a}_{\boldsymbol{\mu}} - e^{-i\mathbf{k}_{\boldsymbol{\mu}} \mathbf{x} + i\omega_{\boldsymbol{\mu}} t} \hat{a}_{\boldsymbol{\mu}}^\dagger]$$

wobei

$$\boldsymbol{\mu} = (\underbrace{\mu_1, \mu_2, \mu_3, \mu_4}_{\mathbf{m}}) \in \mathbb{Z}^3 \times \{\pm 1\} \quad , \quad \mathbf{k}_{\boldsymbol{\mu}} = \mathbf{k}_{\mathbf{m}} = \frac{2\pi}{L} \mathbf{m}$$

$$\mathbf{e}_{\boldsymbol{\mu}} \perp \mathbf{k}_{\boldsymbol{\mu}} \quad , \quad \mathbf{e}_{(\mathbf{m},1)} \perp \mathbf{e}_{(\mathbf{m},2)} \quad , \quad \omega_{\boldsymbol{\mu}} = \omega_{\mathbf{m}} = c \|\mathbf{k}_{\mathbf{m}}\| \quad , \quad \mathbf{e}_{(-\mathbf{m},p)} \stackrel{\text{o.B.d.A.}}{=} \mathbf{e}_{(\mathbf{m},p)}, \quad p = \pm 1$$

und schreiben

$$\begin{aligned}
\hat{\mathbf{p}}_S &= \frac{1}{c^2 \mu_0} \int_V d\mathbf{x} \, \hat{\mathbf{E}} \times \hat{\mathbf{B}} \\
&= -\frac{\hbar}{2V} \int_V d\mathbf{x} \sum_{\mu, \nu \in \mathbb{Z}^3 \times \{\pm 1\}} \mathbf{e}_\mu \times (\mathbf{k}_\nu \times \mathbf{e}_\nu) \cdot \left[e^{i\mathbf{k}_\mu \mathbf{x} - i\omega_\mu t} \hat{a}_\mu - e^{-i\mathbf{k}_\mu \mathbf{x} + i\omega_\mu t} \hat{a}_\mu^\dagger \right] \circ \left[e^{i\mathbf{k}_\nu \mathbf{x} - i\omega_\nu t} \hat{a}_\nu - e^{-i\mathbf{k}_\nu \mathbf{x} + i\omega_\nu t} \hat{a}_\nu^\dagger \right] \\
&= -\frac{\hbar}{2V} \sum_{\mu, \nu \in \mathbb{Z}^3 \times \{\pm 1\}} \mathbf{e}_\mu \times (\mathbf{k}_\nu \times \mathbf{e}_\nu) \cdot \int_V d\mathbf{x} \\
&\quad \left[e^{i(\mathbf{k}_\mu + \mathbf{k}_\nu) \mathbf{x} - i(\omega_\mu + \omega_\nu) t} \hat{a}_\mu \hat{a}_\nu + e^{-i(\mathbf{k}_\mu + \mathbf{k}_\nu) \mathbf{x} + i(\omega_\mu + \omega_\nu) t} \hat{a}_\mu^\dagger \hat{a}_\nu^\dagger - e^{i(\mathbf{k}_\nu - \mathbf{k}_\mu) \mathbf{x} - i(\omega_\nu - \omega_\mu) t} \hat{a}_\mu^\dagger \hat{a}_\nu - e^{i(\mathbf{k}_\mu - \mathbf{k}_\nu) \mathbf{x} - i(\omega_\mu - \omega_\nu) t} \hat{a}_\mu \hat{a}_\nu^\dagger \right] \\
&= -\frac{\hbar}{2} \sum_{\substack{\mathbf{m}, \mathbf{n} \in \mathbb{Z}^3 \\ p, q \in \{\pm 1\}}} \mathbf{e}_{(\mathbf{m}, p)} \times (\mathbf{k}_{\mathbf{n}} \times \mathbf{e}_{(\mathbf{n}, q)}) \cdot \\
&\quad \cdot \left[\delta_{\mathbf{m}, -\mathbf{n}} e^{-i(\omega_{\mathbf{m}} + \omega_{\mathbf{n}}) t} \hat{a}_{(\mathbf{m}, p)} \hat{a}_{(\mathbf{n}, q)} + \delta_{\mathbf{m}, -\mathbf{n}} e^{i(\omega_{\mathbf{m}} + \omega_{\mathbf{n}}) t} \hat{a}_{(\mathbf{m}, p)}^\dagger \hat{a}_{(\mathbf{n}, q)}^\dagger \right. \\
&\quad \left. - \delta_{\mathbf{m}, \mathbf{n}} e^{-i(\omega_{\mathbf{n}} - \omega_{\mathbf{m}}) t} \hat{a}_{(\mathbf{m}, p)}^\dagger \hat{a}_{(\mathbf{n}, q)} - \delta_{\mathbf{m}, \mathbf{n}} e^{-i(\omega_{\mathbf{m}} - \omega_{\mathbf{n}}) t} \hat{a}_{(\mathbf{m}, p)} \hat{a}_{(\mathbf{n}, q)}^\dagger \right] \\
&= -\frac{\hbar}{2} \sum_{\mathbf{m} \in \mathbb{Z}^3} \sum_{p, q \in \{\pm 1\}} \mathbf{e}_{(\mathbf{m}, p)} \times \left[e^{-i2\omega_{\mathbf{m}} t} (\mathbf{k}_{-\mathbf{m}} \times \underbrace{\mathbf{e}_{(-\mathbf{m}, q)}}_{\mathbf{e}_{(\mathbf{m}, q)}}) \cdot \hat{a}_{(\mathbf{m}, p)} \hat{a}_{(-\mathbf{m}, q)} + e^{i2\omega_{\mathbf{m}} t} (\mathbf{k}_{-\mathbf{m}} \times \underbrace{\mathbf{e}_{(-\mathbf{m}, q)}}_{\mathbf{e}_{(\mathbf{m}, q)}}) \cdot \hat{a}_{(\mathbf{m}, p)}^\dagger \hat{a}_{(-\mathbf{m}, q)}^\dagger \right. \\
&\quad \left. - (\mathbf{k}_{\mathbf{m}} \times \mathbf{e}_{(\mathbf{m}, q)}) \cdot \hat{a}_{(\mathbf{m}, p)}^\dagger \hat{a}_{(\mathbf{m}, q)} - (\mathbf{k}_{\mathbf{m}} \times \mathbf{e}_{(\mathbf{m}, q)}) \cdot \hat{a}_{(\mathbf{m}, p)} \hat{a}_{(\mathbf{m}, q)}^\dagger \right] \\
&\stackrel{(\clubsuit)}{=} -\frac{\hbar}{2} \sum_{\substack{\mathbf{m} \in \mathbb{Z}^3 \\ p \in \{\pm 1\}}} \mathbf{k}_{\mathbf{m}} \left[-e^{-i2\omega_{\mathbf{m}} t} \hat{a}_{(\mathbf{m}, p)} \hat{a}_{(-\mathbf{m}, p)} - e^{i2\omega_{\mathbf{m}} t} \hat{a}_{(\mathbf{m}, p)}^\dagger \hat{a}_{(-\mathbf{m}, p)}^\dagger - \hat{a}_{(\mathbf{m}, p)}^\dagger \hat{a}_{(\mathbf{m}, p)} - \hat{a}_{(\mathbf{m}, p)} \hat{a}_{(\mathbf{m}, p)}^\dagger \right]
\end{aligned}$$

wobei verwendet wurde

$$(\clubsuit) : \mathbf{e}_{(\mathbf{m}, p)} \times (\mathbf{k}_{\pm \mathbf{m}} \times \mathbf{e}_{(\mathbf{m}, q)}) = \delta_{p, q} \cdot \mathbf{k}_{\pm \mathbf{m}}$$

Wegen

$$\begin{aligned}
& \sum_{\mathbf{m} \in \mathbb{Z}^3} \mathbf{k}_{\mathbf{m}} \left[e^{-i2\omega_{\mathbf{m}}t} \hat{a}_{(\mathbf{m},p)} \hat{a}_{(-\mathbf{m},p)} + e^{i2\omega_{\mathbf{m}}t} \hat{a}_{(\mathbf{m},p)}^\dagger \hat{a}_{(-\mathbf{m},p)}^\dagger \right] \\
&= \frac{1}{2} \sum_{\mathbf{m} \in \mathbb{Z}^3} \left[\mathbf{k}_{\mu} e^{-i2\omega_{\mathbf{m}}t} \hat{a}_{(\mathbf{m},p)} \hat{a}_{(-\mathbf{m},p)} + \mathbf{k}_{\mu} e^{i2\omega_{\mathbf{m}}t} \hat{a}_{(\mathbf{m},p)}^\dagger \hat{a}_{(-\mathbf{m},p)}^\dagger \right] \\
&+ \frac{1}{2} \sum_{\mathbf{m} \in \mathbb{Z}^3} \left[\underbrace{\mathbf{k}_{\mathbf{m}}}_{-\mathbf{k}_{-\mathbf{m}}} \underbrace{e^{-i2\omega_{\mathbf{m}}t}}_{e^{-i2\omega_{-\mathbf{m}}t}} \underbrace{\hat{a}_{(\mathbf{m},p)} \hat{a}_{(-\mathbf{m},p)}}_{\hat{a}_{(-\mathbf{m},p)} \hat{a}_{(\mathbf{m},p)}} + \underbrace{\mathbf{k}_{\mathbf{m}}}_{-\mathbf{k}_{(-\mathbf{m},p)}} \underbrace{e^{i2\omega_{\mathbf{m}}t}}_{e^{i2\omega_{-\mathbf{m}}t}} \underbrace{\hat{a}_{(\mathbf{m},p)}^\dagger \hat{a}_{(-\mathbf{m},p)}^\dagger}_{\hat{a}_{(-\mathbf{m},p)}^\dagger \hat{a}_{(\mathbf{m},p)}^\dagger} \right] \\
&\stackrel{-\mathbf{m}=\mathbf{n}}{=} \frac{1}{2} \sum_{\mathbf{m} \in \mathbb{Z}^3} \left[\mathbf{k}_{\mathbf{m}} e^{-i2\omega_{\mathbf{m}}t} \hat{a}_{(\mathbf{m},p)} \hat{a}_{(-\mathbf{m},p)} + \mathbf{k}_{\mathbf{m}} e^{i2\omega_{\mathbf{m}}t} \hat{a}_{(\mathbf{m},p)}^\dagger \hat{a}_{(-\mathbf{m},p)}^\dagger \right] \\
&+ \frac{1}{2} \sum_{\mathbf{n} \in \mathbb{Z}^3} \left[-\mathbf{k}_{\mathbf{n}} e^{-i2\omega_{\mathbf{n}}t} \hat{a}_{(\mathbf{n},p)} \hat{a}_{(-\mathbf{n},p)} - \mathbf{k}_{\mathbf{n}} e^{i2\omega_{\mathbf{n}}t} \hat{a}_{(\mathbf{n},p)}^\dagger \hat{a}_{(-\mathbf{n},p)}^\dagger \right] \\
&= 0
\end{aligned}$$

ist sogar

$$\boxed{\hat{\mathbf{p}}_S = \frac{\hbar}{2} \sum_{\boldsymbol{\mu} \in \mathbb{Z}^3 \times \{\pm 1\}} \mathbf{k}_{\boldsymbol{\mu}} \cdot [\hat{a}_{\boldsymbol{\mu}}^\dagger \hat{a}_{\boldsymbol{\mu}} + \hat{a}_{\boldsymbol{\mu}} \hat{a}_{\boldsymbol{\mu}}^\dagger]} \quad (0.3)$$

Somit ergibt sich

$$\hat{\mathbf{p}}_S |n_1, n_2, \dots\rangle = \frac{\hbar}{2} \sum_{\boldsymbol{\mu}} \mathbf{k}_{\boldsymbol{\mu}} [n_{\boldsymbol{\mu}} + (n_{\boldsymbol{\mu}} + 1)] |n_1, n_2, \dots\rangle$$

wobei

$$\sum_{\boldsymbol{\mu} \in \mathbb{Z}^3 \times \{\pm 1\}} \mathbf{k}_{\boldsymbol{\mu}} = \frac{1}{2} \sum_{\boldsymbol{\mu}} \left[\mathbf{k}_{\boldsymbol{\mu}} + \underbrace{\mathbf{k}_{\boldsymbol{\mu}}}_{-\mathbf{k}_{-\boldsymbol{\mu}}} \right] \stackrel{-\boldsymbol{\mu} \rightarrow \boldsymbol{\mu}}{=} \frac{1}{2} \sum_{\boldsymbol{\mu}} [\mathbf{k}_{\boldsymbol{\mu}} - \mathbf{k}_{\boldsymbol{\mu}}] = 0$$

also

$$\boxed{\hat{\mathbf{p}}_S |n_1, n_2, \dots\rangle = \sum_{\boldsymbol{\mu} \in \mathbb{Z}^3 \times \{\pm 1\}} \hbar n_{\boldsymbol{\mu}} \mathbf{k}_{\boldsymbol{\mu}} |n_1, n_2, \dots\rangle} \quad (0.4)$$

□