

Quantenfeldtheorie

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Aufgabe 25

(a) Wegen $\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu}$ antivertauscht γ^μ mit allen $\gamma^{\nu \neq \mu}$. Daher:

$$\gamma_5 \gamma^\mu = i\gamma^0 \gamma^1 \gamma^2 \gamma^3 \gamma^\mu = (-1)^3 i\gamma^\mu \gamma^0 \gamma^1 \gamma^2 \gamma^3 = -\gamma^\mu \gamma_5$$

(b) Aus $\gamma^0 \gamma^0 = 1$, $\gamma^i \gamma^i = -1$ folgt

$$\gamma_5^2 = -\gamma^0 \gamma^1 \gamma^2 \gamma^3 \gamma^0 \gamma^1 \gamma^2 \gamma^3 = -(-1)^3 (-1)^2 (-1)^1 (-1)^0 \underbrace{\gamma^0 \gamma^0}_1 \underbrace{\gamma^1 \gamma^1}_{-1} \underbrace{\gamma^2 \gamma^2}_{-1} \underbrace{\gamma^3 \gamma^3}_{-1} = 1$$

(c)

$$P_{R,L}^2 = \frac{1}{4}(1 \pm 2\gamma_5 + \underbrace{\gamma_5^2}_1) = \frac{1}{2}(1 \pm \gamma_5) = R_{R,L}$$

$$P_R P_L = \frac{1}{4}(1 + \cancel{\gamma_5} - \cancel{\gamma_5} - \underbrace{\gamma_5^2}_1) = 0$$

$$P_R + P_L \stackrel{\text{klar}}{=} 1$$

(d) Die Lagrange-Dichte des Dirac-Felder ist gegeben durch

$$\mathcal{L} = i\bar{\psi}\gamma^\mu \partial_\mu \psi - m\bar{\psi}\psi, \quad \bar{\psi} := \psi^\dagger \gamma^0$$

Dabei gilt

$$\bar{\psi}_L \stackrel{\text{def.}}{=} (P_L \psi)^\dagger \gamma^0 = \psi^\dagger P_L^\dagger \gamma^0 = \frac{1}{2}(\psi^\dagger \gamma^0 - \psi^\dagger \underbrace{\gamma_5 \gamma^0}_{-\gamma^0 \gamma_5}) = \frac{1}{2}(\bar{\psi} + \bar{\psi} \gamma_5) = \bar{\psi} P_R$$

$$\bar{\psi}_R \stackrel{\text{analog}}{=} \bar{\psi} P_L$$

Somit kann der kinetische Term geschrieben werden als:

$$\begin{aligned} i\bar{\psi}\gamma^\mu \partial_\mu \psi &= i\bar{\psi} \overbrace{(P_L + P_R)}^1 \gamma^\mu \partial_\mu \overbrace{(P_L + P_R)}^1 \psi \\ &= i\bar{\psi} P_L \gamma^\mu \partial_\mu P_L \psi + i\underbrace{\bar{\psi} P_L}_{\bar{\psi}_R} \gamma^\mu \partial_\mu \underbrace{P_R \psi}_{\psi_R} + i\underbrace{\bar{\psi} P_R}_{\bar{\psi}_L} \gamma^\mu \partial_\mu \underbrace{P_L \psi}_{\psi_L} + i\bar{\psi} P_R \gamma^\mu \partial_\mu P_R \psi \\ &= i\bar{\psi}_R \gamma^\mu \partial_\mu \psi_R + i\bar{\psi}_L \gamma^\mu \partial_\mu \psi_L + i\bar{\psi} \underbrace{(P_L \gamma^\mu P_L + P_R \gamma^\mu P_R)}_{\substack{P_R \gamma^\mu \\ \gamma^\mu P_L}} \partial_\mu \psi \end{aligned}$$

$$P_L P_R = 0 \quad i\bar{\psi}_R \gamma^\mu \partial_\mu \psi_R + i\bar{\psi}_L \gamma^\mu \partial_\mu \psi_L$$

Andererseits ist

$$\overline{\psi}_L \psi_L + \overline{\psi}_R \psi_R = \overline{\psi} \underbrace{P_R P_L}_0 \psi + \overline{\psi} \underbrace{P_L P_R}_0 \psi$$

das heißt eine ähnliche Separation ist für den Massenterm nicht möglich!

□

Aufgabe 26

Die Transformation $\psi \mapsto e^{i\alpha} \psi$ erhält offensichtlich den Lagrangian:

$$\mathcal{L}(\psi) \mapsto \mathcal{L}(e^{i\alpha} \psi) = e^{-i\alpha} \overline{\psi} i \partial_\mu \gamma^\mu e^{i\alpha} \psi - \frac{\lambda}{2} \left[(e^{-i\alpha} \overline{\psi} e^{i\alpha} \psi)^2 - (e^{-i\alpha} \overline{\psi} \gamma_5 e^{i\alpha} \psi)^2 \right] = \mathcal{L}(\psi)$$

was den Noether-Strom

$$j^\nu = \underbrace{\frac{\partial \mathcal{L}}{\partial(\partial_\nu \psi)}}_{-i\overline{\psi} \gamma^\nu} (i\alpha) \psi + (i\alpha) \overline{\psi} \underbrace{\frac{\partial \mathcal{L}}{\partial(\partial_\nu \overline{\psi})}}_0 = -\alpha \overline{\psi} \gamma^\nu \psi$$

liefert. Wegen

$$\exp[-i\alpha \gamma_5] \gamma^\mu = \sum_{n=0}^{\infty} \frac{(-i\alpha)^n}{n!} \underbrace{\gamma_5^n \gamma^\mu}_{(-1)^n \gamma^\mu \gamma_5^n} = \gamma^\mu \exp[i\alpha \gamma_5]$$

führt $\psi \mapsto e^{i\alpha \gamma_5} \psi$ auf den transformierten Lagrangian

$$\begin{aligned} \mathcal{L}(\psi) \mapsto \mathcal{L}(e^{i\alpha \gamma_5} \psi) &= i\overline{\psi} \underbrace{e^{i\alpha \gamma_5} \gamma^\mu}_{\gamma^\mu e^{-i\alpha \gamma_5}} \partial_\mu e^{i\alpha \gamma_5} \psi - \frac{\lambda}{2} \left[\underbrace{(\overline{\psi} e^{i\alpha \gamma_5} e^{i\alpha \gamma_5} \psi)^2}_{e^{2i\alpha \gamma_5}} - \underbrace{(\overline{\psi} e^{i\alpha \gamma_5} \gamma_5 e^{i\alpha \gamma_5} \psi)^2}_{\gamma_5 e^{i\alpha \gamma_5}} \right] \\ &= i\overline{\psi} \gamma^\mu \partial_\mu \psi - \frac{\lambda}{2} \left[(\overline{\psi} e^{2i\alpha \gamma_5} \psi)^2 - (\overline{\psi} \gamma_5 e^{2i\alpha \gamma_5} \psi)^2 \right] \end{aligned}$$

Die Wirkung des Generators $i\alpha \gamma_5$ auf \mathcal{L} ergibt sich also gemäß

$$\begin{aligned} \Delta_{i\alpha \gamma_5} \mathcal{L} &= \partial_t \mathcal{L}(e^{it\alpha \gamma_5} \psi) \Big|_{t=0} = \alpha \cdot [\partial_\alpha \mathcal{L}(e^{i\alpha \gamma_5} \psi)]_{\alpha=0} = -\lambda \alpha [(\overline{\psi} \psi) (\overline{\psi} 2i\gamma_5 \psi) - (\overline{\psi} \gamma_5 \psi) (\overline{\psi} \gamma_5 2i\gamma_5 \psi)] \\ &= -2i\lambda \alpha \underbrace{[(\overline{\psi} \psi) (\overline{\psi} \gamma_5 \psi) - (\overline{\psi} \gamma_5 \psi) (\overline{\psi} \psi)]}_{[\overline{\psi} \psi, \overline{\psi} \gamma_5 \psi] = 0} = 0 \end{aligned}$$

Dementsprechend erhält man den Noether-Strom

$$j^\mu = \underbrace{\frac{\partial \mathcal{L}}{\partial(\partial_\nu \psi)}}_{-i\overline{\psi} \gamma^\nu} \underbrace{\Delta_{i\alpha \gamma_5}}_{i\alpha \gamma_5 \psi} + \Delta_{i\alpha \gamma_5} \overline{\psi} \underbrace{\frac{\partial \mathcal{L}}{\partial(\partial_\nu \overline{\psi})}}_0 = \alpha \overline{\psi} \gamma^\nu \gamma_5 \psi$$