

Quantenfeldtheorie

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Aufgabe 07

Schreiben

$$\begin{aligned} -a_\mu \{\phi(x), P^\mu\} &= -a_\mu \int d^d x \left[\underbrace{\delta_{\phi(z)} \phi(x)}_{\delta(x-z)} \cdot \delta_{\pi(z)} P^\mu - \underbrace{\delta_{\pi(z)} \phi(x)}_0 \cdot \delta_{\phi(z)} P^\mu \right] \\ &= -a_\mu \delta_{\pi(x)} P^\mu = -a_\mu \int d^d z \left[\underbrace{(\delta_{\pi(x)} \pi(z))}_{\delta(z-x)} \cdot \partial^\mu \phi(z) + \pi(z) \underbrace{\delta_{\pi(x)} \partial^\mu \phi(z)}_{\partial^\mu \delta_{\pi(x)} \phi(z)=0} - g^{\mu 0} \delta_{\pi(x)} \mathcal{L} \right] \\ &= -a_\mu \partial^\mu \phi(x) + a^\mu \int d^d z g^{\mu 0} \left[\frac{\partial \mathcal{L}}{\partial \phi} \underbrace{\delta_{\pi(x)} \phi(z)}_0 + \frac{\partial \mathcal{L}}{\partial (\partial_\nu \phi)} \underbrace{\delta_{\pi(x)} \partial_\nu \phi(z)}_{\partial_\nu \delta_{\pi(x)} \phi(z)=0} \right] \\ &= -a_\mu \partial^\mu \phi(x) = -a^\mu \partial_\mu \phi(x) = \delta \phi(x) \end{aligned}$$

Aufgabe 08

Kanonische Impulsdichte

Mit

$$\phi = \frac{1}{\sqrt{2}} [\phi_1 + i\phi_2]$$

und

$$\mathcal{L} = \sum_{k=1}^2 \left\{ \frac{1}{2} (\partial_\mu \phi_k) (\partial^\mu \phi_k) - \frac{m^2}{2} \phi_k^2 \right\}$$

ergibt sich

$$\begin{aligned} \pi_k &= \frac{\partial \mathcal{L}}{\partial \dot{\phi}_k} = \frac{1}{2} \left[\underbrace{\frac{\partial (\partial_\mu \phi_k)}{\partial (\partial_0 \phi_k)}}_{\delta_{0\mu}} (\partial_\nu \phi_k) g^{\mu\nu} + (\partial_\mu \phi_k) \underbrace{\frac{\partial (\partial_\nu \phi_k)}{\partial (\partial_0 \phi_k)}}_{\delta_{0\nu}} g^{\mu\nu} \right] = \frac{1}{2} [(\partial_\nu \phi_k) g^{0\nu} + (\partial_\mu \phi_k) g^{\mu 0}] \\ &= (\partial_\nu \phi_k) g^{0\nu} = \partial^0 \phi_k \end{aligned}$$

und somit

$$\partial^0 \phi = \frac{1}{\sqrt{2}} [\pi_1 + i\pi_2] =: \pi \quad .$$

Hamilton-Dichte

Schreiben

$$\begin{aligned}\mathcal{H} &= \sum_{k=1}^2 \pi_k \dot{\phi}_k - \mathcal{L} = \sum_{k=1}^2 \left[\frac{1}{2} \pi_k^2 + \frac{1}{2} \nabla \phi_k \cdot \nabla \phi_k + \frac{m^2}{2} \phi_k^2 \right] \\ &= \frac{1}{2} (\pi_1 - i\pi_2) (\pi_1 + i\pi_2) + \frac{1}{2} (\nabla \phi_1 - i\nabla \phi_2) (\nabla \phi_1 + i\nabla \phi_2) + \frac{m^2}{2} (\phi_1 - i\phi_2) (\phi_1 + i\phi_2) \\ &= |\pi|^2 + \nabla \phi^* \cdot \nabla \phi + \frac{m^2}{2} |\phi|^2\end{aligned}$$

Spektrale Darstellung der Hamilton-Funktion

Nennen $\mathcal{F}(f)(k) =: \tilde{f}(k)$ die Fourier-Transformation von f und schreiben

$$\begin{aligned}H &= \int d^d x \mathcal{H}(\phi(x), \pi(x)) = \int d^d x \left[|\pi|^2 + \sum_{m=1}^d \partial_m \phi^* \cdot \partial_m \phi + m^2 |\phi|^2 \right] \\ &\stackrel{\text{Parseval}}{=} \int d^d k \left[|\tilde{\pi}(k)|^2 + m^2 |\tilde{\phi}(k)|^2 \right] + \frac{1}{(2\pi)^n} \sum_{m=1}^d \int d^d x \int d^d k \underbrace{\mathcal{F}(\partial_m \phi^*)(k)}_{\substack{ik_m \mathcal{F}(\phi^*)(k) \\ = ik_m (\mathcal{F}(\phi))^*(-k)}} \cdot e^{ikx} \cdot \int d^d k' \underbrace{\mathcal{F}(\partial_m \phi)(k')}_{{ik'_m \mathcal{F}(\phi)}} \cdot e^{ik'x} \\ &= \int d^d k \left[|\tilde{\pi}(k)|^2 + m^2 |\tilde{\phi}(k)|^2 \right] - \frac{1}{(2\pi)^n} \sum_{m=1}^d \int \int d^d k d^d k' k_m k'_m \cdot (\mathcal{F}(\phi))^*(-k) \cdot \mathcal{F}(\phi)(k') \cdot \underbrace{\int d^d x e^{i(k+k')x}}_{(2\pi)^n \delta(k+k')} \\ &= \int d^d k \left[|\tilde{\pi}(k)|^2 + m^2 |\tilde{\phi}(k)|^2 \right] + \sum_{m=1}^d \int d^d k' k_m^2 (\mathcal{F}(\phi))^*(k') \mathcal{F}(\phi)(k') \\ &= \int d^d k \left[|\mathcal{F}(\pi)(k)|^2 + \|k\|^2 |\mathcal{F}(\phi)(k)|^2 + m^2 |\mathcal{F}(\phi)(k)|^2 \right]\end{aligned}$$

Quantisierung der Feldoperatoren

Beginnend mit

$$[a_k(\mathbf{p}), a_j(\mathbf{q})] = [a_k^\dagger(\mathbf{p}), a_j^\dagger(\mathbf{q})] = 0 \quad , \quad [a_k(\mathbf{p}), a_j^\dagger(\mathbf{q})] = \underbrace{(2\pi)^d}_{\varkappa} \delta_{kj} \delta(\mathbf{p} - \mathbf{q})$$

schreiben wir

$$[a(\mathbf{p}), a(\mathbf{q})] = \frac{1}{2} [a_1(\mathbf{p}) + ia_2(\mathbf{p}), a_2(\mathbf{q}) + ia_1(\mathbf{q})] = \underbrace{[a_1(\mathbf{p}), a_1(\mathbf{q})]}_0 + i \underbrace{[a_1(\mathbf{p}), a_2(\mathbf{q})]}_0 + i \underbrace{[a_2(\mathbf{p}), a_1(\mathbf{q})]}_0 - \underbrace{[a_2(\mathbf{p}), a_2(\mathbf{q})]}_0 = 0$$

$$[b(\mathbf{p}), b(\mathbf{q})] \stackrel{\text{Analog}}{=} 0 \quad , \quad [a(\mathbf{p}), b(\mathbf{q})] = -[a^\dagger(\mathbf{p}), b^\dagger(\mathbf{q})] \stackrel{\text{Analog}}{=} 0$$

$$[a(\mathbf{p}), a^\dagger(\mathbf{q})] = \frac{1}{2} \underbrace{[a_1(\mathbf{p}), a_1^\dagger(\mathbf{q})]}_{\varkappa \delta(\mathbf{p}-\mathbf{q})} + \frac{1}{2} \underbrace{[a_2(\mathbf{p}), a_2^\dagger(\mathbf{q})]}_{\varkappa \delta(\mathbf{p}-\mathbf{q})} = \varkappa \delta(\mathbf{p} - \mathbf{q}) \quad , \quad [b(\mathbf{p}), b^\dagger(\mathbf{q})] \stackrel{\text{Analog}}{=} \varkappa \delta(\mathbf{p} - \mathbf{q})$$

$$[a(\mathbf{p}), b^\dagger(\mathbf{q})] = \frac{1}{2} \underbrace{[a_1(\mathbf{p}), a_1^\dagger(\mathbf{q})]}_{\varkappa \delta(\mathbf{p}-\mathbf{q})} - \frac{1}{2} \underbrace{[a_2(\mathbf{p}), a_2^\dagger(\mathbf{q})]}_{\varkappa \delta(\mathbf{p}-\mathbf{q})} = 0$$

Mit

$$a_1(\mathbf{p}) = \frac{1}{\sqrt{2}} [a(\mathbf{p}) + b(\mathbf{p})] \quad , \quad a_1^\dagger(\mathbf{p}) = \frac{1}{\sqrt{2}} [a^\dagger(\mathbf{p}) + b^\dagger(\mathbf{p})]$$

$$a_2(\mathbf{p}) = \frac{i}{\sqrt{2}} [b(\mathbf{p}) - a(\mathbf{p})] \quad , \quad a_2^\dagger(\mathbf{p}) = \frac{i}{\sqrt{2}} [a^\dagger(\mathbf{p}) - b^\dagger(\mathbf{p})]$$

und

$$\phi_k(\mathbf{p}) = \frac{1}{\sqrt{2\omega_{\mathbf{p}}}} [a_k(\mathbf{p}) + a_k^\dagger(-\mathbf{p})] \quad , \quad \pi_k(\mathbf{p}) = -i\sqrt{\frac{\omega_{\mathbf{p}}}{2}} [a_k(\mathbf{p}) - a_k^\dagger(-\mathbf{p})]$$

ergibt sich ferner

$$\phi(\mathbf{p}) = \frac{1}{\sqrt{2}} [\phi_1(\mathbf{p}) + i\phi_2(\mathbf{p})] = \frac{1}{\sqrt{2\omega_{\mathbf{p}}}} [a(\mathbf{p}) + b^\dagger(-\mathbf{p})]$$

$$\pi(\mathbf{p}) = \frac{1}{\sqrt{2}} [\pi_1(\mathbf{p}) + i\pi_2(\mathbf{p})] = i\sqrt{\frac{\omega_{\mathbf{p}}}{2}} [b^\dagger(-\mathbf{p}) - a(\mathbf{p})]$$

Darstellung des Hamilton-Operators

Eingesetzt in obiger Darstellung für H ergibt

$$\begin{aligned} H &= \frac{1}{\alpha} \int d^d \mathbf{p} [\pi(\mathbf{p})\pi^\dagger(\mathbf{p}) + \omega_{\mathbf{p}}^2 \cdot \phi(\mathbf{p})\phi^\dagger(\mathbf{p})] \\ &= \frac{1}{\alpha} \int d^d \mathbf{p} \omega_{\mathbf{p}} \cdot [\underbrace{a(\mathbf{p})a^\dagger(\mathbf{p})}_{\varkappa(\delta(0)+a^\dagger(\mathbf{p})a(\mathbf{p}))} + b^\dagger(-\mathbf{p})b(-\mathbf{p})] \\ &= \frac{1}{(2\pi)^d} \int d^d \mathbf{p} \omega_{\mathbf{p}} \cdot [a^\dagger(\mathbf{p})a(\mathbf{p}) + b^\dagger(\mathbf{p})b(\mathbf{p})] + \underbrace{\int d^d \mathbf{p} \omega_{\mathbf{p}} \delta(0)}_{\text{N.P.E.}} \end{aligned}$$

Noether-Ladung

Schreiben Q um:

$$\begin{aligned} Q &= i \int d^d \mathbf{x} [\phi^\dagger(\mathbf{x})\pi(\mathbf{x}) - \phi(\mathbf{x})\pi^\dagger(\mathbf{x})] \\ &= \frac{i}{\varkappa^2} \int \int \int d^d \mathbf{x} d^d \mathbf{p} d^d \mathbf{q} [\phi^\dagger(\mathbf{p})\pi(\mathbf{q})e^{i\mathbf{x}(\mathbf{q}-\mathbf{p})} - \phi(\mathbf{p})\pi^\dagger(\mathbf{q})e^{i\mathbf{x}(\mathbf{p}-\mathbf{q})}] \\ &= \frac{i}{\varkappa^2} \int d^d \mathbf{p} d^d \mathbf{q} \left[\phi^\dagger(\mathbf{p})\pi(\mathbf{q}) \underbrace{\int d^d \mathbf{x} e^{i\mathbf{x}(\mathbf{q}-\mathbf{p})}}_{\varkappa\delta(\mathbf{q}-\mathbf{p})} - \phi(\mathbf{p})\pi^\dagger(\mathbf{q}) \underbrace{\int d^d \mathbf{x} e^{i\mathbf{x}(\mathbf{p}-\mathbf{q})}}_{\varkappa\delta(\mathbf{p}-\mathbf{q})} \right] \\ &= \frac{i}{(2\pi)^d} \int d^d \mathbf{p} [\phi^\dagger(\mathbf{p})\pi(\mathbf{p}) - \phi(\mathbf{p})\pi^\dagger(\mathbf{p})] \end{aligned}$$

Durch obige Darstellungen von ϕ und π ferner

$$\begin{aligned} Q &= \frac{i}{\varkappa} \int d^d \mathbf{p} \frac{i}{2} \left\{ [a^\dagger(\mathbf{p}) + b(-\mathbf{p})] \cdot [b^\dagger(-\mathbf{p}) - a(\mathbf{p})] + [a(\mathbf{p}) + b^\dagger(-\mathbf{p})] \cdot [b(-\mathbf{p}) - a^\dagger(\mathbf{p})] \right\} \\ &= \frac{1}{2\varkappa} \int d^d \mathbf{p} \left[\underbrace{a(\mathbf{p})a^\dagger(\mathbf{p})}_{a^\dagger a + \varkappa\delta(0)} + a^\dagger(\mathbf{p})a(\mathbf{p}) - b^\dagger(-\mathbf{p})b(-\mathbf{p}) - \underbrace{b(-\mathbf{p})b^\dagger(-\mathbf{p})}_{b^\dagger b + \varkappa\delta(0)} \right] \\ &= \frac{1}{\varkappa} \int d^d \mathbf{p} [a^\dagger(\mathbf{p})a(\mathbf{p}) - b^\dagger(-\mathbf{p})b(-\mathbf{p})] = \frac{1}{(2\pi)^d} \int d^d \mathbf{p} [a^\dagger(\mathbf{p})a(\mathbf{p}) - b^\dagger(\mathbf{p})b(\mathbf{p})] \end{aligned}$$